# **A surface energy analysis of bioadhesion\***

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This report applies recently developed surface energy and fracture mechanics relations to the analysis of bioadhesion and biocompatibility. The dispersion  $\alpha$  and polar  $\beta$  components of 190 biological and implant surfaces are analysed. The surface energetics relations between bioadhesion and biocompatibility point out that a strongly adsorbed plasma protein film on the implant surface provides the best blood compatibility and low thrombogenic effects. The surface energy relations provide means of selecting optimum implant surface properties and mapping on surface energy diagrams the three phase interactions which define bioadhesion.

# INTRODUCTION

A number of literature reports document the general progress in development of biomaterials and surface treatments designed for application in cardiovascular devices<sup> $1-8$ </sup>. The interfacial interaction of biopolymers with blood remains as one of the central problems in establishing biocompatibility of cardiovascular devices<sup>8</sup>. The exposure of blood to a foreign surface produces a complex set of concurrent and sequential events which appear to correlate with the dispersion (London-d) and polar (Keesom-p) components of surface tension for the implant material. Baier and coworkers $1-7$  have reported detailed studies of contact angle measurements of well characterized liquids on biopolymer surfaces. These measurements and surface energetics analysis follow the methodology and definitions of critical surface tension  $\gamma_c$  for wetting of a solid substrate developed by Zisman and coworkers<sup>9</sup>.

More recently, Nyilas and coworkers<sup>10</sup>, have demonstrated that the wettability data of Baler and coworkers can be utilized in defining the dispersion and polar components of solid surface tension for candidate implant surfaces. This report of Nyilas and coworkers also describes new semiquantitative relations between blood flow, implant surface energetics, and thrombosis. The analysis utilized by Nyilas and coworkers to isolate the dispersion and polar properties of implant surfaces follows definitions and calculations developed and extensively applied by Kaelble and cowor $kers<sup>11-13</sup>$ . The surface energetics analysis of Kaelble and coworkers has recently been extended<sup>14</sup> to define the relations between surface energetics and the Griffith fracture mechanics criteria for spontaneous interface bonding and debonding under consitions of combined liquid immersion and added mechanical stress.

This paper discusses the results of applying the new surface energetics criterion of bonding and debonding for improving the quantitative definition of bioadhesion and to clarify the relationship between bioadhesion and biocompatibility. These new surface energy relations now permit mapping the zones of bonding, termed wettability envelopes on surface energy diagrams of dispersion  $\alpha$  versus polar  $\beta$ components of surface energy. These surface energy *a versus*   $\beta$  diagrams permit graphic presentation of the surface properties and zones of bonding and debonding for the three phases which constitute the interfacial boundary in the biocompatibility analysis.

### SURFACE ENERGY ANALYSIS

The general concept for regular adsorption bonding of interfaces utilized in this discussion is summarized in the following relation for interfacial tension<sup>12</sup>:

$$
\gamma_{ij} = (\alpha_i - \alpha_j)^2 + (\beta_i - \beta_j)^2 + \Delta_{ij}
$$
 (1)

where the parameters are defined in *Table 1* and subscripts denote interactions from phase *i* and *j*. Interfaces dominated by Van der Waal's interactions are termed regular interfaces

*Table 1* Surface energetics relations

$\gamma_{LV} = \gamma_{l} \frac{d}{V} + \gamma_{l} \frac{p}{V} = \alpha_L^2 + \beta_L^2$	(a)
$\gamma_{SV}$ = $\gamma_{SV}^d$ + $\gamma_{SV}^p$ = $\alpha_S^2$ + $\beta_S^2$	(b)
$W_a = \gamma_{LV} (1 + \cos \theta) \leq 2\gamma_{LV}$	(c)
$W_a = 2[\alpha_L \alpha_S + \beta_L \beta_S]$	(d)
$\frac{W_a}{\sqrt{2}} = \alpha_S + \beta_S(\beta_L/\alpha_L)$ $2\alpha_L$	(e)

$$
\frac{W_3}{2\alpha_S} = \alpha_L + \beta_L (\beta_S/\alpha_S) \tag{f}
$$

where:  $\gamma_{LV}$  = liquid-vapour tension;  $\gamma_{SV}$  = solid-vapour surface tension;  $\alpha_1, \beta_1$  = square root of the respective (London) dispersion  $\gamma_L V$  and (Keesom) polar  $\gamma_L V$  parts of  $\gamma_L$   $_V$ ;  $\alpha_S$   $\beta_S$  = square roots of respective dispersion  $\gamma\mathcal{S}_{\mathcal{V}}$  and polar  $\gamma\mathcal{S}_{\mathcal{V}}$ ;  $W_{\mathcal{A}}$  = nominal work of adhesion;  $\theta$  = liquid—solid contact angle

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*Table 2* Fracture mechanics relations

$$
\sigma_C = \left(\frac{2E\gamma_G}{\pi_C}\right)^{1/2} = \left(\frac{2E}{\pi C}\right)^{1/2} \left(\frac{2E}{\pi C}\right)^{1/2} = 0
$$
 (a)

$$
\gamma_G = R^2 - R_0^2 \tag{b}
$$

 $R_0^2 = 0.25 \left[ (\alpha_1 - \alpha_3)^2 + (\beta_1 - \beta_3)^2 \right]$  (c)

$$
R^{2} = (\alpha_{2} - H)^{2} + (\beta_{1} - K)^{2}
$$
 (d)

 $H = 0.5 (\alpha_1 + \alpha_3)$  (e)

 $K = 0.5 (\beta_1 + \beta_3)$  (f)

where  $\sigma_c$  = critical crack propagation stress;  $\gamma_G$  = Griffith surface energy for fracture;  $E = \text{Young's modulus}; c = \text{crack length}; \alpha_1, \beta_1 =$ surface properties of adhesive (ink) phase 1;  $\alpha_2,\alpha_2$  = surface properties of environment phase  $2$ ;  $\alpha_3$ ,  $\beta_3$  = surface properties of adherend phase 3

and the values of the excess term  $\Delta_{ij}$  of equation (1) which describes interdiffusion or ionic-covalent interactions can be considered negligible. This is a much more general case than one might expect and permits application of surface energy analysis to a wide range of materials. When  $\Delta_{ii} = 0$ , equation (1) defines an ideal interface with  $\gamma_{ij}$  = o as the special case where  $\alpha_i = \alpha_j$  and  $\beta_i = \beta_j$ .

The special combination of surface energy and fracture mechanics parameters which enter the modified Griffith relation are defined in *Table 2* and show that the Griffith fracture energy  $\gamma_G$  is defined by the following relation:

$$
\gamma_G = -\frac{S_2}{2} = R^2 - R_0^2 \tag{2}
$$

A circular parabola in  $\gamma_G$ ,  $\alpha_2$ ,  $\beta_2$  Cartesian space is defined by equation (2). The surface energies  $\alpha_2$  and  $\beta_2$  for the immersion phase 2 which provide the condition  $R < R_0$  provides that the spreading coefficient  $S_2$  for phase 2 is positive with  $S_2 > 0$ . The predicted consequence for  $S_2 > 0$  is that phase 2 should spontaneously debond phase 1 from phase 3 in the absence of rheological restraints. When  $R \leq R_0$  the Griffith fracture energy becomes positive and a critical mechanical stress  $\sigma_c$  which depends on  $\gamma_G$  (see *Table 2*) is now required for crack extension.

The relations of *Table I and Table 2* form the basis for designed experiments which isolate the discrete mechanisms of polar and dispersion interactions across interface. The test liquids of *Table 3* display a wide range of polar character in surface tension with  $\beta_l/\alpha_l = 1.53$  for water to  $\beta_l/\alpha_l =$ 0 for linear hydrocarbons. Inspection of equation (e) in *Table 1* shows that using measured values of work of adhesion *Wa* by contact angle measurements for liquids of known  $\alpha_l$  and  $\beta_l$  permits isolation of the solid-vapour surface properties  $\alpha_s$  and  $\beta_s$ . The intercept of the plot of  $W_a/2\alpha_l$  versus  $\beta_l/\alpha_l$  isolates  $\alpha_s$  as the intercept and  $\beta_s$  as the slope. An alternative method of computerized determinant solutions of equation (d) (Table 1) has been described and extensively applied which solve for averaged values of  $\gamma_{s\nu}^d$ ,  $\gamma_{s\nu}^p$  and  $\gamma_{s\nu}$ their respective standard deviations  $\pm \delta d$ ,  $\delta p$ , and  $\delta$  from mean values $^{11,12}$ .

In this study contact angle data where  $\theta > 0$  and  $W_a < 2\gamma_{lv}$ published by Baier and coworkers $1-7$  is combined with the liquid surface tension data of *Table 3 in* the computerized determinant calculations for  $\gamma_{sv}^a$ ,  $\gamma_{sv}^p$ ,  $\gamma_{sv}$  and  $\pm \delta^a$ ,  $\delta^p$ , and  $\delta$ . The surface energetics of 190 biological and implant surfaces

were analysed and the results summarized as values of  $\alpha_s$  =  $(\gamma_{sy}d)^{1/2}$  and  $\beta_s = (\gamma_{sy}d)^{1/2}$  and the percent standard deviation  $(\delta \times 100/\gamma_{s\nu})$  of  $\gamma_{s\nu}$  in *Table 4*. The surface number sequence of *Table 4* correlates with the appearance of the original experimental data in the referenced literature to facilitate ease of cross reference. The reference articles and reports cover an important six year period of biomaterials development and testing from 1970 through 1975.

Surfaces no. 1 and no. 2 of *Table 4* represent important biological surfaces and the illustrative results of the surface energy analysis are graphed in *Figure 1. Figure lb* shows the wettability data for human fibrinogen thin film as defined by equation (e) of *Table 1*. The solid linear curve of *Figure lb* graphs the computer calculated average values  $\gamma_{\rm w}^d$  = 24.6 and  $\gamma_{\rm w}^p$  = 13.5 dyne/cm while the broken curves define the standard deviation  $\delta^{d} = \pm 0.7$  and  $\delta^{b} = \pm 0.9$  dyne/cm. As shown in *Figure 1b* the experimental values of  $W_a/2\alpha_l$  form a reasonably linear curve which conforms well with the theory of *Table i* and the computer solutions.

The data points and linear curves of *Figure la* show the larger uncertainty in  $\alpha_s$  and  $\beta_s$  values for vein intimal surface no. 2 of *Table 4.* Referring to the discussion of Baier *et al.*  it is evident that this surface is soft and deformable, the experiment complicated, additionally possible interdiffusion modifies the contact angle data for water and glycerol. As shown in *Figure la,* this analysis permits isolation of a low dispersion and high polar surface energy for the vein intimal surface. Considering the highly hydrated state of the vein intimal surface, it is to be expected that the vein surface properties  $\alpha = 4.42$  (dyne/cm)<sup>1/2</sup> and  $\beta = 6.18$  (dyne/cm)<sup>1/2</sup> are closely similar to the surface tension properties of water where  $\alpha$  = 4.67 (dyne/cm)<sup>1/2</sup> and  $\beta$  = 7.14 (dyne/cm)<sup>1/2</sup>.

Available literature references<sup>14,15</sup> describe the surface tension of blood plasma as equivalent to saline solution with surface tension  $\gamma_{lv}$  = 72 to 74 dyne/cm, which is essentially equivalent to pure water with  $\gamma_{lv}$  = 72.8 dyne/cm. Further studies may show that surfactant effects of blood plasma constituents can produce variable values of  $\beta_l$  with nearly constant  $\alpha_l$  in which  $\beta_l = 7.14$  (dyne/cm)<sup>1/2</sup> represents a maximum value. This result has been described by Kaelble<sup>16</sup> for aqueous detergent solutions above the critical micelle concentration. For this discussion, the surface properties of pure water with  $\alpha_l$  = 4.67 and  $\beta_l$  = 7.14 (dyne/cm)<sup>1/2</sup> are

*Table 3* Surface **tension properties of test liquids** at 20°C

<b>Test liquid</b>	~Iv (dyne/cm)	$\alpha$	βι	Reference
Water	72.8	4.67	7.14	10, 13
Glycerol	63.4	6.10	5.12	10
Formamide	58.2	6.28	4.32	10
Dithioglycol	54.0	6.20	3.94	10
Methylene iodide	50.8	6.83	2.05	10
Ethylene glycol	48.3	5.41	4.36	13
S-Tetrabromoethane	47.5	6.49	2.32	10
$\alpha$ -Bromonaphthalene	44.6	6.68	3.59	10
Tricresyl phosphate	40.9	6.26	1.30	13
1-Methylnaphthalene	38.7	5.04	3.65	10
Dicyclohexyl	33.0	5.74	0	10
n-Hexadecane	27.6	5.24	O	13
n-Tetradecane	26.7	5.17	0	10
n-Trídecane	25.9	5.09	0	10
n-Dodecane	25.4	5.04	0	10
n-Decane	23.9	4.89	0	10
n-Nonane	22.8	4.77	0	10
n-Octane	21.8	4.67	0	10
n-Heptane	20.3	4.51	0	10











 $\mathcal{L}^{\text{max}}_{\text{max}}$ 

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taken as the analogue for the surface properties of blood plasma. This assumption sets forth the proposition that protein molecular segments (phase 1) are competing with the water molecules of blood plasma (phase 2) for bonding sites on the implant surface (phase 3).

The surface energy properties of fibrinogen, vein intima, can be graphically displayed on a surface energy diagram of  $\alpha$  versus  $\beta$  as shown in *Figure 2*. This diagram also locates the surface tension properties of water ( $\approx$ blood plasma) and graphically displays the close relation between the vein intima surface properties and those of water. Inserting the appropriate values for  $\alpha$  and  $\beta$  into equation (1) for  $\Delta_{ij}$  = 0 one calculates an interfacial tension  $\gamma_{ij}$  = 12.1 dyne/cm between blood and fibrinogen as compared to a much lower value of  $\gamma_{ii}$  = 1.0 dyne/cm between the vein intima surface and water. The vein surface thus provides a much closer approach to an ideal interface where  $\gamma_{ij} = 0$  than does the fibrinogen Film.

## BIOADHESION AND BIOCOMPATIBILITY

Substantial evidence $1-8$  now shows that the first event that follows the exposure of blood to an implant material is the adsorption of plasma borne protein which covers the implant

surface. This adsorbed plasma protein film should modify the surface energy of the implant so as to lower its interfacial tension to blood plasma. From equation (1) the specific set of interfacial tensions between the implant (phase 3), blood plasma (phase 2) and adsorbed plasma protein film (phase 1) are started as follows:

$$
\gamma_{12} = (\alpha_1 - \alpha_2)^2 + (\beta_1 - \beta_2)^2 \tag{3}
$$

$$
\gamma_{13} = (\alpha_1 - \alpha_3)^2 + (\beta_1 - \beta_3)^2 \tag{4}
$$

$$
\gamma_{23} = (\alpha_2 - \alpha_3)^2 + (\beta_2 - \beta_3)^2 \tag{5}
$$

For adsorbed plasma protein (phase 1) to spontaneously displace blood plasma (phase 2) from the implant (phase 3) the spreading coefficient S, for phase 1 must be positive. The phase 1 spreading coefficient is defined as:

$$
S_1 = \gamma_{23} - \gamma_{13} - \gamma_{12} \tag{6}
$$

As shown in equation (6), protein adsorption is favoured by large values of  $\gamma_{23}$ . As shown in equation (5) a large mismatch in implant surface properties  $\alpha_3$ ,  $\beta_3$  and blood plasma  $\alpha_2 \approx 4.67, \beta_2 \approx 7.14 \text{ (dyne/cm)}^{1/2}$  is seen to increase  $\gamma_{23}$ 



*Figure I* Analytic definition of surface tension properties of (a) vein intimal surface, dog(surface 2 of *Table 4*),  $yd = 19.5 \pm 2.6$ ;  $\gamma p = 38.2 \pm 9.5$  dyne cm;  $\alpha = 4.42$ ;  $\beta = 6.18$ ; (b) human fibrinogen thin film (surface 1 of *Table 4*),  $\gamma d = 2.46 \pm 0.7$ ;  $\gamma p = 13.5 \pm 0.9$ dyne/cm,  $\alpha$  = 4.96;  $\beta$  = 3.67

and enhance protein adsorption.

The complementary proposition in equation (6) is that minimum values of both  $\gamma_{13}$  and  $\gamma_{12}$  also favour protein adsorption. Equation (3) shows that a minimum mismatch in the surface properties  $\alpha_1 \approx \alpha_2$  and  $\beta_1 \approx \beta_2$  reduces  $\gamma_{12}$ . A similar minimization of  $\gamma_{13}$  is obtained by achieving  $\alpha_1 \approx \alpha_3$ and  $\beta_1 \approx \beta_{13}$  as shown in equation (4). The present analysis of protein analogues, see surfaces no. 22-23 in *Table 4,*  shows that the polar  $\beta$  part of surface energy is substantially varied by solvent-polymer interactions and resultant polymer chain conformation. It is reasonable to presume that plasma protein will be adsorbed on an implant surface with chain conformations which tend to minimize both  $\gamma_{12}$  and  $\gamma_{13}$  by spontaneous adjustment of  $\alpha_1$  and  $\beta_1$  values in the protein film. By spontaneously minimizing  $\gamma_{12}$  and  $\gamma_{13}$ , the adsorbed plasma protein film evidently performs an important biological adaptation function.

The physiochemical description of the plasma displacement and protein adsorption on the implant surface is analogous to 'priming' as familiarly described in paint technology. Attachment of platelets with formation of thrombus and subsequent thrombus release to produce embolism appears to be mediated by the protein prime layer. Clinical tests of biocompatibility are briefly described in the Appendix by two in viva tests in common usage. The vena cava test developed by Gott and coworkers<sup>17,19</sup> is limited to the detection of thrombus. The renal embolus test, developed by Kusserow and coworkers<sup>18,19</sup> detects both thrombus and embolism generated by thrombus release from the implant surface. Recent studies by Baier and coworkers<sup>7</sup> now reveal that implant surfaces with evident high thrombo-

resistance in the vena cava test are shown to be thrombogenic in the renal embolus test where inspection of the kidney reveals extensive implant damage. The indications are that the thrombus can form on the implant surface and continuously spall off to be carried and deposited in the kidney. The end result is a relatively thrombus free implant surface that acts as a continuous embolus generator. Simple inspection of the implant surfaces at the site of implantation is therefore not a sufficient test for biocompatibility.

A central issue in a revised defmition of biocompatibility can be related to the resistance to detachment of the plasma protein film which 'primes' and biologically modifies the implant surface. The modified Griffith relations of *Table 2*  provide a quantitative means for evaluating the Griffith energy  $\gamma_G$  required to detach the adsorbed protein film (phase 1) from the implant (phase 3) in the presence of blood plasma (phase 2). The studies summarized in *Table 4*  include surface energy analysis of implant materials before and after implantation. Surfaces no. 14-21 studied by Baier and coworkers<sup>2</sup> furnish data for the calculations of the Griffith surface energy  $\gamma$ G as summarized in *Table 5.* 

The calculations in *Table 5* show that the Griffith surface energy  $\gamma_G$  and related critical mechanical stress  $\sigma_G$  decrease as the polar component  $\beta_3$  of the implant surface increases. The clinical ratings of biocompatibility listed in the lower portion of *Table 5* show a direct correlation between increased  $\gamma_G$  and improved blood compatibility. The results of *Table 5* can be mapped on surface energy diagrams of  $\alpha$ *versus*  $\beta$  as shown in *Figure 3. Figure 3b* for polished carbon defines high  $\gamma_G$  with decreasing values of  $\gamma_G$  in clockwise direction to show surface cleaned stellite 21 metal in the lower right view as most highly thrombogenic. The surface energy diagrams of *Figure 3* show the points H, K defined in *Table 2* as the origin of the R and  $R_0$  vectors which define the Griffith fracture energy:

$$
\gamma_G = R^2 - R_0^2 \tag{7}
$$

as defined in *Table 2*. The magnitude of  $R_0$  which subtracts



*Figure 2* Dispersion  $(\alpha)$  and polar  $(\beta)$  surface properties of biolo**gical** materials





\* TMDAC = tridodecylammonium chloride; *GDT* = glow discharge treated

from  $\gamma_G$  is related to the mismatch in surface properties between implant (phase 3) and adsorbed protein layer (phase 1). Conversely the manitude of R which adds to  $\gamma_G$  is related to the mismatch between  $H$ ,  $K$  which defines an averaged property of the 1-3 interface and the blood plasma phase. As shown in *Table 5* in all cases the magnitude of R2 dominantly controls the magnitude of  $\gamma_G$  in the examples graphed in *Figure 3.* The design concept for blood compatibility described by equation (7) incorporates a general argument which details the competition between blood plasma (phase 2) and plasma protein (phase 1) for bonding sites on the implant surface. A high  $\gamma$ <sub>G</sub>, indicative of blood compatibility, favours both the spontaneous formation and strong retention of the adsorbed protein film.

#### DISCUSSION

The previous two sections have briefly introduced and illustrated methods for analysis of surface energy and bioadhesion. The detailed discussion of all relevant aspects of the extensive data compilation in *Table 4* is beyond the scope of this brief report. Review of *Table 4* in conjunction with the extensive experimentation and discussion by Baier and coworkers in the original references<sup> $1-7$ </sup> shows general agreement with the results illustrated here wherein a combination of high dispersion ( $\alpha$ ) combined with low ( $\beta$ ) correlates with high blood compatibility. The previous section relates this result with the strong adsorption and stable retention of a plasma protein adsorption layer on the implant surface.

If the virgin implant surface has a highly polar character which approaches the  $\beta$  values for water or blood plasma the examples show that the protein layer is weakly adsorbed and held on the implant surface. In this later case, the incompletely covered implant would appear to operate efficiently as a thrombus and embolus generator.

This preliminary application of surface energy analysis and the modified Griffith criteria of *Table 1* and *Table 2,*  respectively, is quite encouraging. Nyilas and coworkers<sup>10</sup> have reported a series of studies of thrombus formation under closely controlled blood flow conditions. Within a given category of implant materials a high polar property  $\beta$ for the implant surface is shown to correlate with low thromboresistance for a given shear rate of blood flow. This result of Nyilas and coworkers<sup>10</sup> is in agreement with the modified Griffith analysis of critical scress  $\sigma_c$  or energy  $\gamma_G$  for debonding applied in this report.

The adsorption theory outlined in *Table 1* and *Table 2*  and applied in this report does not directly treat the effects of either surface roughness or interdiffusion effects related to solvolytic interactions at the interface. Kaelble<sup>12</sup> has developed the thermodynamic extensions of adsorption theory of interfaces to treat both roughness and interdiffusion. The effects of microroughness at the implant surface is potentially a dominant issue in surface energetics at the microfibre scaffold surface used for anchoring viable fibroplastic and endothelial cells to produce a 'living' blood compatible surface.

The adsorption theory is also ill equipped to deal quantitatively with time dependent changes in bioadhesion arising from interdiffusion effects. This latter point is graphically evident in the data scatter shown for blood vessel intima in *Figure la.* The natural linings of blood vessels are hydrogels. Synthetic hydrogel coatings consisting of a coherent three dimensional polymer network containing a large proportion of water display promise as implant surfaces<sup>8</sup>. A more detailed description of the role of both *adsorbed* and *absorbed* water on the bioadhesion of hydrogel coatings and biologically deposited protein films require use of a combined adsorption interdiffusion (A-I) theory.

The right column of *Table 4* lists the percent standard deviation from the mean  $(100 \delta \gamma_{s\nu}^{-1})$  for computed average values of  $\gamma_{s\nu}$ . Of the 190 surfaces examined, 134 display standard deviations of less than 5% and 175 surfaces give deviations less than 10%. The maximum standard deviation for solid No. 14 is 21.47%. Large standard deviations are generally related to roughness and interdiffusion effects and readily identified in data displays as shown in *Figure 1.*  Imprecise values of  $\alpha$  and  $\beta$  are, of course, carried forward into the calculations of  $\gamma$ <sub>G</sub> as presented in *Table 5.* 



*Figure 3* Surface energy analysis of interracial interactions between plasma protein (phase 1), blood plasma  $\approx$  water (phase 2), and implant (phase 3). (a) 1, 3 (above) after implantation; 2, water  $\approx$  blood plasma; 3, diamond polished carbon. (b) 1, 3 (above) after  $implantation; 2$ , water  $\approx$  blood plasma; 3, TDMAC heparanized silicone. (c) 1, 3 (above after implantation; 2, water  $\approx$  blood plasma; 3, GDT Stellite 21. (d) 1, 3 (above) after implantation; 2, water blood plasma; 3, polyurethane electret

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#### **CONCLUSIONS**

As a result of this study, some conclusions concerning the role of surface energetics in bioadhesion can be made as follows:

(A) Surface energy analysis is successfully applied to define the dispersion  $(\alpha)$  and polar  $(\beta)$  surface energies of 190 biological and implant surfaces.

(B) These dispersion-polar surface energies are introduced into fracture mechanics relations for bioadhesion and biocompatibility.

(C) These fracture mechanics calculations indicate that blood compatibility of an implant is enhanced by spontaneous absorption and strong retention of a plasma protein film on the implant surface.

(D) High dispersion-low polar surface energy for the implant as exemplified by low temperature isotropic  $(LTI)$ carbon with  $\alpha \ge 6.0$  (dyne/cm)<sup>1/2</sup>, and  $\beta \le 2.0$  (dyne/cm)<sup>1/2</sup>, provide surface energetics favouring stable plasma protein film retention.

(E) Low dispersion-high polar surfaces, typified by surface treated Stellite 21 with  $\alpha \approx 5.0$  (dyne/cm)<sup>1/2</sup>,  $\beta \geq 5.0$  (dyne/cm)<sup>1/2</sup>, provide surface energetics, appear to favour weak adsorption and retention of the plasma protein such that the implant may continuously generate and spall off emboli into the blood stream.

(F) The present analysis can be extended to describe the effects of interface roughness and interdiffusion which represent dominant considerations in microfibre scaffold surfaces, hydrogel coatings, and biological intimal surfaces of the cardiovascular system.

The extensive listing of surface energies in *Table 4* is intended to be used in conjunction with the referenced data sources. The main objective in the discussion is to illustrate, by examples, the usefulness of this surface energy analysis (see *Figure 1)* and the application of the analysis of spontaneous bonding and debonding (see *Figures 2* and 3).

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# APPENDIX

#### *Clinical tests for inplant bioeompatibility*

*Vena cava test.* This test was developed by Gott and coworkers<sup>17,19</sup>. Rings (length = 9 mm, and diameters  $o.d.$  = 8 mm and i.d. 7 mm) with streamlined edges to prevent turbulence are fabricated from the test material or surface coated with the test material. Implantation is made in the inferior vena cava of a 7.7 to 19.4 dg dog. The rings are inserted with a special device to insure noncontact with the atrial wall or contamination with tissue fluid during implantation.

For *acute* studies (2 h) the chest remains open. At the end of the two hour period the vena cava above and below the ring is doubly clamped and the ring is quickly excised. The inside of the ring is examined within a minute. The extent and nature of the gross thrombus is recorded immediately on a standard ring chart.

For *chronic* studies (2 weeks) the same implantation techniques are used. After implantation, the chest is closed and the animal maintained two weeks before sacrifice. The chest cavity is quickly opened, the ring excised and examined for the extent and nature of thrombus and results recorded on the standard ring chart. In addition, there is a gross inspection of the lungs for any pulmonary pathology and pulmonary embolae.

*Renal embolus test.* This test was developed by Kusserow and coworkers<sup>18,19</sup>. A ring (length  $= 10$  mm, and diameters  $o.d. = 8.6$  mm and i.d. = 7 mm) is implanted into the abdominal aorta of the test animal immediately above the origin renal arteries. A construction is made in the aorta slightly below the origin of the renal arteries so that over 90% of the blood flowing through the ring must pass through the kidneys. After an implantation period of three to five days an autopsy is performed. The extent of surface thrombus of the ring is assessed by direct visual observation. Embolism is evaluated by direct visual and microscopic examination of both kidneys. The kidneys serve as efficient biological accumulators for the embolic phenomena because the emboli produce recognizable infarcts in these organs.